

TSALLIS-THERMOMETER: A QGP INDICATOR FOR LARGE AND SMALL COLLISIONAL SYSTEMS

20TH ZIMÁNYI SCHOOL WINTER WORKSHOP ON HEAVY ION PHYSICS

GÁBOR BÍRÓ

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WIGNER RESEARCH
CENTRE FOR **PHYSICS**
EÖTVÖS LORÁND **UNIVERSITY**

Collaborators:

GERGELY GÁBOR BARNAFÖLDI
TAMÁS SÁNDOR BIRÓ

Talk based on:

G. Bíró, G.G. Barnaföldi, T.S. Biró, J. Phys. G, 47.10 (2020), 105002.

Related publications:

G. Bíró, G.G. Barnaföldi, K. Ürmösy, T.S. Biró, Á. Takács, Entropy, 19(3), (2017), 88

G. Bíró, G.G. Barnaföldi, T.S. Biró, K. Shen, EPJ Web Conf, 171, (2018), 14008

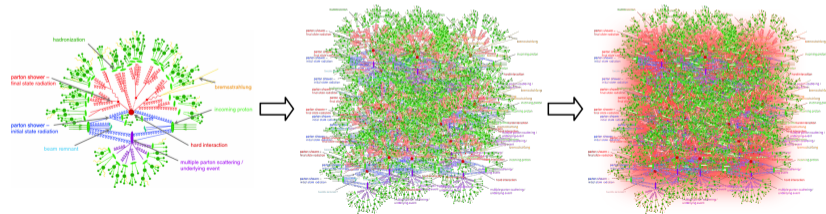
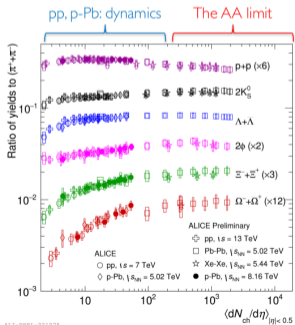
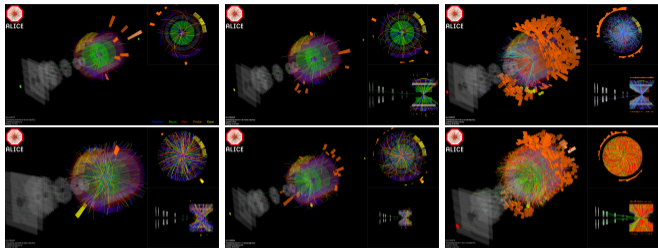


MOTIVATION

Ratio of identified hadrons in small to large systems...

...but what is **small**?

Small systems can have **large** multiplicities too...



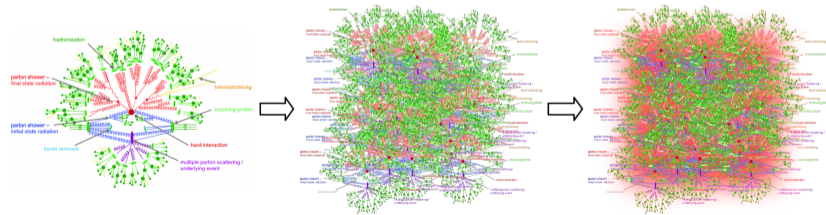
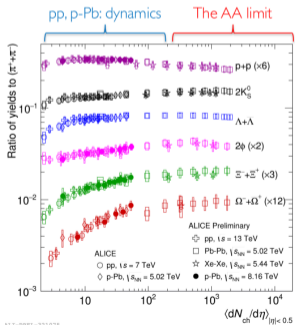
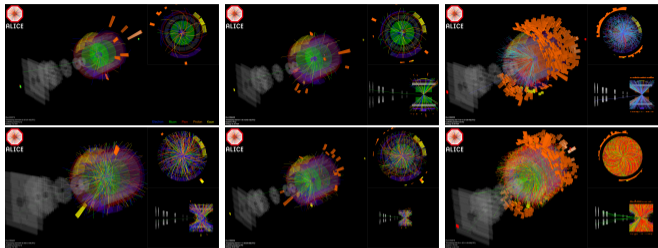
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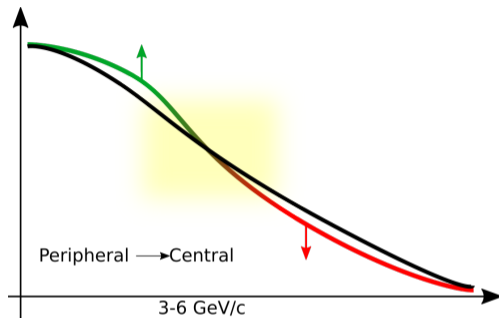
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Small systems can have **large** multiplicities too...

Where does the quark-gluon plasma start in **multiplicity**?



Non-extensive statistics – summary:



$$\frac{d^2N}{2\pi p_T dp_T dy} = A m_T \left[1 + \frac{q-1}{T} (m_T - m) \right]^{-\frac{q}{q-1}}$$

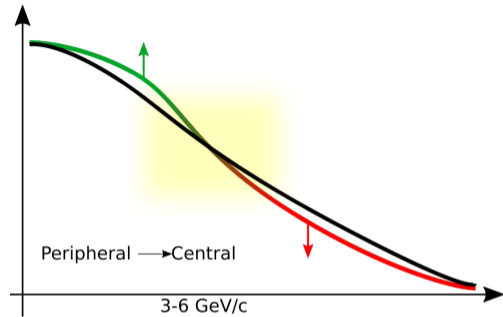
Entropy 16(12), (2014), 6497-651. **Eur.Phys.J.A** 55 (2019) 8, 126

Non-extensive statistics – summary:

q -entropy:

$$S_q = \frac{1}{q-1} \left(1 - \sum_{i=1}^W p_i^q \right)$$

$$\lim_{q \rightarrow 1} S_q = S_{BG}$$



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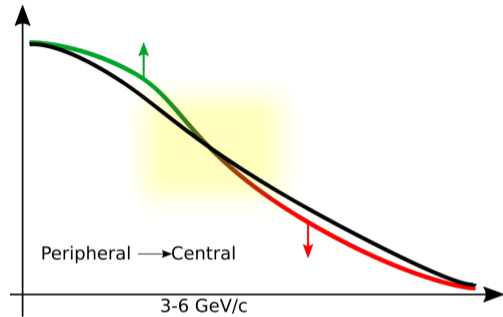
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Thermodynamical consistency:

$$P = Ts + \mu n - \varepsilon$$

$$P = g \int \frac{d^3 p}{(2\pi)^3} T f \qquad s = g \int \frac{d^3 p}{(2\pi)^3} \left[\frac{E - \mu}{T} f^q + f \right]$$

$$N = nV = gV \int \frac{d^3 p}{(2\pi)^3} f^q \qquad \varepsilon = g \int \frac{d^3 p}{(2\pi)^3} E f^q$$



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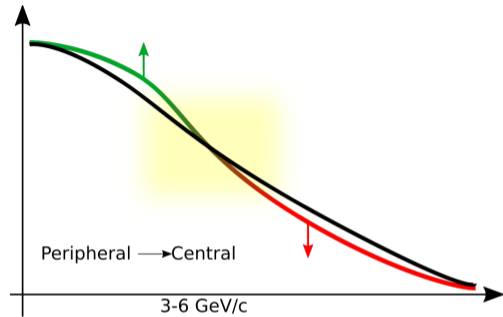
Final size effects:

$$T = \frac{E}{\langle n \rangle}$$

$$q = 1 - \frac{1}{\langle n \rangle} + \frac{\Delta n^2}{\langle n \rangle^2}$$

$$T = E \left[\delta^2 - (q-1) \right]$$

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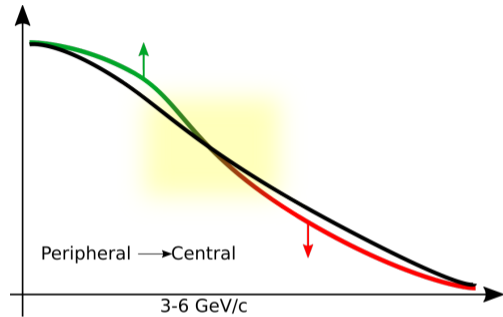
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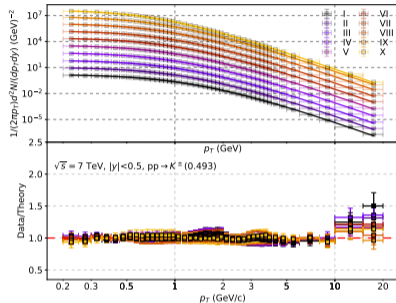
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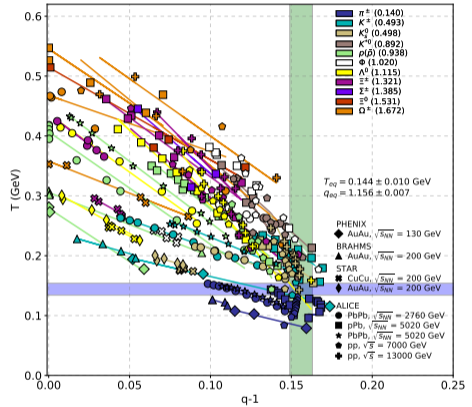
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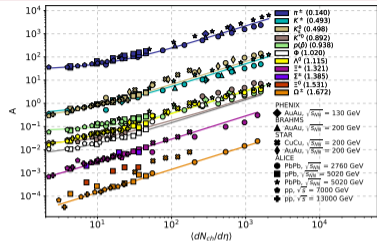
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Phenomenological approach:

Map the thermodynamically consistent non-extensive parameter space of the available experimental data and compare it with theoretical QCD calculations

- 11 identified hadron species: from π^\pm to Ω
- Various collision systems: proton-proton, proton-nucleus, nucleus-nucleus
- Wide range of multiplicities: $2.2 \leq \langle dN_{ch}/d\eta \rangle \leq 2047$
- Wide range of CM energies: $130 \leq \sqrt{s_{NN}} \leq 13000$ GeV
- **More than 30** published experimental datasets



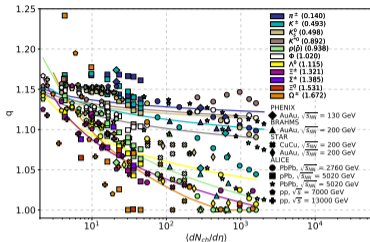
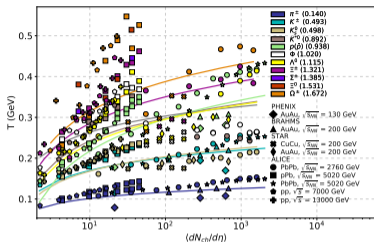


Parametrizations:

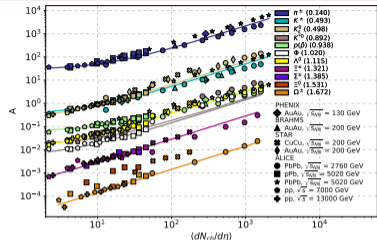
$$A = A_0 + A_1 \ln \frac{\sqrt{s_{NN}}}{m} + A_2 \langle dN_{ch}/d\eta \rangle$$

$$T = T_0 + T_1 \ln \frac{\sqrt{s_{NN}}}{m} + T_2 \ln \ln \langle dN_{ch}/d\eta \rangle$$

$$q = q_0 + q_1 \ln \frac{\sqrt{s_{NN}}}{m} + q_2 \ln \ln \langle dN_{ch}/d\eta \rangle$$



1. The **A**, **q** and **T** parameters characterize the collision
2. Strong **grouping**: $T_{eq} \approx 0.144$ GeV, $q_{eq} \approx 1.156$



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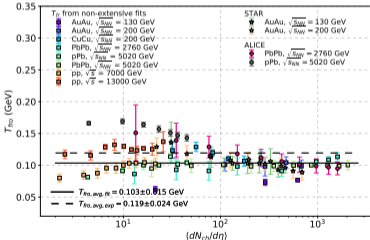
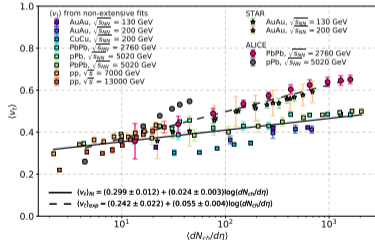
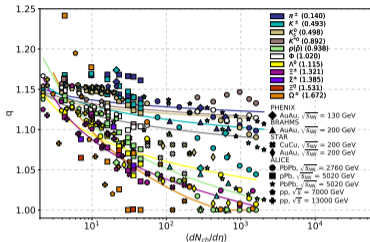
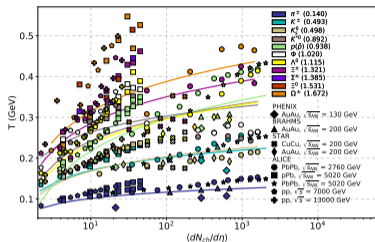
$$T = T_0 + T_1 \ln \frac{\sqrt{s_{NN}}}{m} + T_2 \ln \ln \langle dN_{ch}/d\eta \rangle$$

$$q = q_0 + q_1 \ln \frac{\sqrt{s_{NN}}}{m} + q_2 \ln \ln \langle dN_{ch}/d\eta \rangle$$

Radial flow:

$$T = T_{fro} + m \langle u_t \rangle^2$$

$$\langle v_t \rangle = \frac{\langle u_t \rangle}{\sqrt{1 + \langle u_t \rangle^2}}$$



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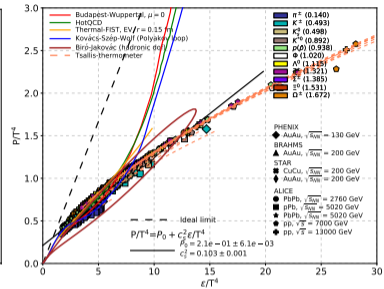
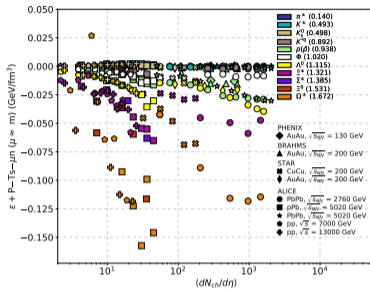
2. Strong **grouping**: $T_{eq} \approx 0.144$ GeV, $q_{eq} \approx 1.156$

3. **Test**: results are comparable with experiments (**Phys. Rev. C 83 (2011), 064903**)

Thermodynamical consistency: ✓

$$P = Ts + \mu n - \varepsilon$$

Comparison of the thermodynamical variables with theoretical calculations

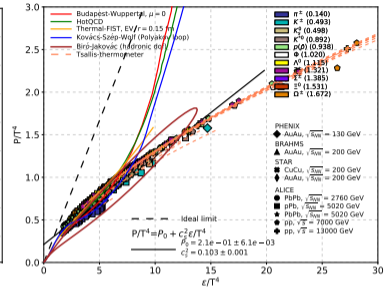
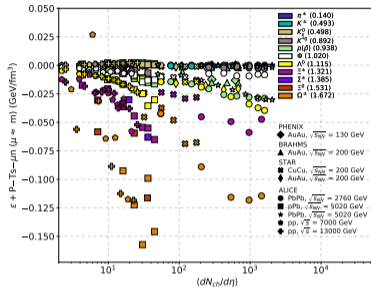


Interpretation of the grouping phenomenon in the $T - (q - 1)$ parameter space:

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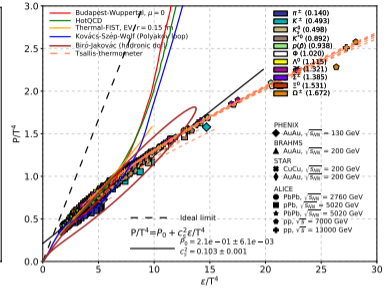
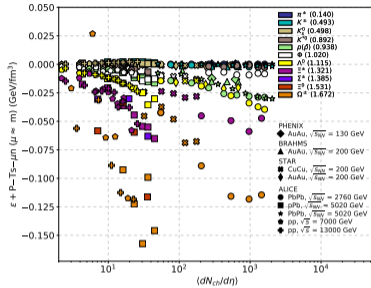
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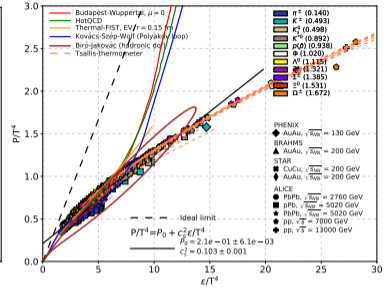
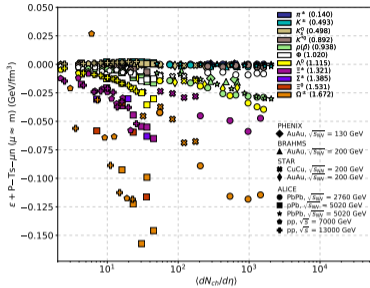
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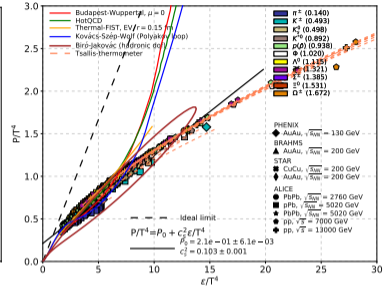
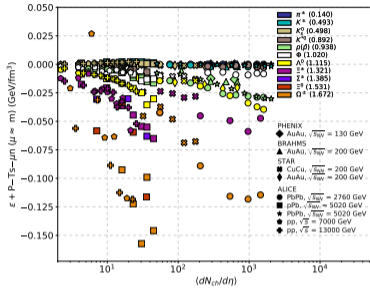
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3. This QGP does certainly **not** follow an equilibrium Boltzmann - Gibbs statistics

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With the **parametrizations**: \sqrt{s} and $\langle dN_{ch}/d\eta \rangle$ regions:

- $\sqrt{s} \gtrsim 7000$ GeV: $\langle dN_{ch}/d\eta \rangle \gtrsim 130$
- $\sqrt{s} \gtrsim 13000$ GeV: $\langle dN_{ch}/d\eta \rangle \gtrsim 90$

SUMMARY

- Consistent non-extensive analysis of a **very large set** of experimental data
- $q \neq 1$ for all hadron spectra: dependency on the size of the collisional system through **multiplicity** fluctuations
- **Various checks** of the non-extensive framework
- Grouping of the **T** and **q** parameters, **comparison** with theoretical QCD calculations
- **Tsallis-thermometer**: final state hadrons may originate from a previously present strongly interacting QCD matter at event multiplicities as low as $\langle dN_{ch}/d\eta \rangle \sim 100$

SUPPORT

The research is supported by: OTKA K120660, K123815, K135515, THOR COST CA15213, Hungarian-Chinese 12 CN-1-2012-0016, MOST 2014DFG02050, 2019-2.1.11-TÉT-2019-00050 Tét, Wigner HAS-OBOR-CCNU, ÚNKP-17-3.

Thank you for your attention!

BACKUP

EXPERIMENTAL DATA

System, $\sqrt{s_{NN}}$ (GeV)	η or y	Hadron	Mult. classes	p_T range (GeV/ c)	
AuAu, 130	$ \eta < 0,35$	π^\pm	3, [21,3; 622]	[0,25; 2,2]	
		K^\pm		[0,45; 1,65]	
CuCu, 200	$ y < 0,5$	$p(\bar{p})$	5, [32; 175]	[0,55; 3,42]	
		K_s^0		[0,5; 9,0]	
		Λ^0		[0,5; 7,0]	
		Ξ^\pm		[0,7; 6,0]	
		Ω^\pm		[1,0; 4,5]	
		Φ		[0,45; 4,5]	
		π^\pm		6, [24; 175]	[0,2; 2,0]
AuAu, 200	$ y < 0,2$	K^\pm	3, [111; 680]	[0,4; 2,0]	
		$p(\bar{p})$		[0,3; 3,0]	
PbPb, 2760	$ y < 0,5$	K_s^0	5, [27; 680]	[0,5; 9,0]	
		Λ^0		[0,5; 8,0]	
		π^\pm		10, [13,4; 1601]	[0,1; 3,0]
		K^\pm		[0,2; 3,0]	
		K_s^0		7, [55; 1601]	[0,4; 12,0]
		K^{*0}		6, [261; 1601]	[0,3; 20,0]
		$p(\bar{p})$		[0,3; 4,6]	
		Λ^0		[0,6; 12,0]	
		Φ		[0,5; 21,0]	
		Ξ^\pm		5, [55; 1601]	[0,6; 8,0]
pPb, 5020	$-0,5 < y < 0,0$	Ω^\pm	7, [4,3; 45]	[1,2; 7,0]	
		π^\pm		[0,1; 20,0]	
		K^\pm		[0,2; 20,0]	
		K^{*0}		5, [4,3; 45]	[0,0; 16,0]
		$p(\bar{p})$		[0,35; 20,0]	
		Φ		[0,4; 20,0]	
		Ξ^0		4, [7,1; 35,6]	[0,8; 8,0]

System, $\sqrt{s_{NN}}$ (GeV)	η or y	Hadron	Mult. classes	p_T range (GeV/ c)				
pPb, 5020	$-0,5 < y < 0,0$	π^\pm	7, [4,3; 45]	[0,1; 20,0]				
		Σ^\pm		3, [7,1; 35,6]	[1,0; 6,0]			
		Ξ^\pm		7, [4,3; 45]	[0,6; 7,2]			
	$0,0 < y < 0,5$	Ω^\pm	7, [4,3; 45]	[0,8; 5,0]				
		π^\pm		[0,1; 3,0]				
		K^\pm		[0,2; 2,4]				
PbPb, 5020	$ y < 0,5$	K_s^0	10, [19,5; 2047]	[0,0; 8,0]				
		$p(\bar{p})$		[0,3; 4,0]				
		Λ^0		[0,6; 8,0]				
		π^\pm		[0,1; 10,0]				
		K^\pm		[0,1; 10,0]				
		$p(\bar{p})$		[0,1; 10,0]				
		pp, 7000		$ y < 0,5$	π^\pm	10, [2,2; 21,3]	[0,1; 20,0]	
					K^\pm		[0,2; 20,0]	
					K_s^0		10, [2,2; 21,3]	[0,0; 12,0]
					K^{*0}		9, [2,2; 21,3]	[0,0; 10,0]
$p(\bar{p})$	10, [2,2; 21,3]		[0,3; 20,0]					
Φ	9, [2,2; 21,3]		[0,4; 10,0]					
Λ^0	10, [2,2; 21,3]		[0,4; 8,0]					
Ξ^\pm	[0,6; 6,5]							
pp, 13000	$ y < 0,5$	Ω^\pm	5, [2,2; 21,3]	[0,9; 5,5]				
		K_s^0	10, [2,52; 25,72]	[0,0; 12,0]				
		Λ^0		[0,4; 8,0]				
		Ξ^\pm		[0,6; 6,5]				
		Ω^\pm		5, [3,58; 22,8]	[0,9; 5,5]			